

Center for Detectors

Avalanche Photodiodes

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Review of Solid State Physics



Band Theory

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Figure 14 Calculated band structure of germanium, after C. Y. Fong. The general features are in good agreement with experiment. The four valence bands are shown in gray. The fine structure of the valence band edge is caused by spin-orbit splitting. The energy gap is indirect; the conduction band edge is at the point $(2\pi/a)(\frac{1}{2}\frac{1}{2})$. The constant energy surfaces around this point are ellipsoidal.







Band Theory



Figure 17b Band structure of GaAs, after S. G. Louie.





The PN Junction

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PN Junction Band Diagram (Zeghbroeck 2007)

I-V Characteristics of a Diode

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Breakdown and the Avalanche Mechanism

 E_{c} E_{c} EFn E_{Fn} E_v E_{ν} (a) Small reverse bias (b) $|V_A| \rightarrow V_{BR}$

Figure 6.12 Carrier activity inside the depletion region of a reverse-biased *pn* junction diode when (a) $|V_A| \ll V_{BR}$ and (b) $|V_A| \rightarrow V_{BR}$. Carrier multiplication due to impact ionization and the resultant avalanche is pictured in (b).



Avalanche Photodiodes



Linear Mode vs Geiger Mode

Once an avalanche has begun, a competition develops between the rate at which electron-hole pairs are generated and the rate at which they are collected at the device terminals.

Linear mode:

At biases below but near the breakdown voltage, collection dominates, causing the avalanche current to decay and ultimately stop. In this mode, the gain is finite and predictable (to within some degree of uncertainty based on material properties). Because the gain is finite, it is possible to distinguish between a single photon or multiple photons arriving simultaneously, though distinguishing between large numbers of photons (i.e., > 5) is statistically degenerate since the gain per electron is probabilistic.

Linear mode APDs can be used as single photon counting devices (low flux levels) or as photodiodes with a built-in gain at higher flux levels.



Linear Mode vs Geiger Mode

(Once an avalanche has begun, a competition develops between the rate at which electron-hole pairs are generated and the rate at which they are collected at the device terminals.)

Geiger mode:

In the case where the bias is above V_{BR} , multiplication outpaces collection. Initially, this causes exponential growth of the current. After some length of time, electrons and holes accumulate at the n- and p-sides of the depletion region, respectively, creating an internal electric field in opposition to the applied bias (arresting the growth of the current). The device remains in an on state, however, until the circuit reduces the applied bias (quenches the device), allowing the APD to turn off. In this mode, the gain would be infinite if the bias were held above breakdown. It is important to note that in this mode, an electrical event resulting from a single incident photon is indistinguishable from one initiated by a larger number of photons arriving simultaneously.

Practical Design

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The guard ring implant (heavily doped region) reduces edge breakdown effects (when the edge of the device has a lower magnitude breakdown voltage than the center).

The cathode is biased such that the voltage over the PN junction is greater than the breakdown voltage.

Practical Design

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Photon incident from back side

The relatively small multiplier region (< 1 µm) facilitates a higher electric field across at lower voltages. This allows for higher gain at lower voltages than the standard structure APD. This particular device is designed for Geiger-mode operation and breaks down at about 28 V reverse bias.



Performance Metrics (LM-APD)

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Leakage Current (un-multiplied dark current)

Not all dark carriers are multiplied. Some travel to the cathode without going through the multiplication region. These carriers will add to the integrated charge, introducing a constant slope to the ideal step function.



Performance Metrics (LM-APD)

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- Dark Count Rate
 - Dark current carriers can be multiplied just like photo-generated carriers.
 - Practically, it is impossible to distinguish between the photo-generated signal and the dark (thermal) carrier noise. Therefore, the number of expected dark counts is subtracted from the final signal (though the variation from the noise will remain, decreasing SNR).





Performance Metrics (LM-APDs)

Excess Noise Factor

 The ratio of electron and hole impact ionization rates (k) affects the uncertainty in the gain per avalanche. This is a material property. The excess noise factor (F) is a function of the mean gain and k:

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F=kM+(1-k)(2-M^{-1})
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Therefore, even the ideal shot-noise limited case is affected:

 $SNR = S/sqrt(S^*F)$

- Silicon has a k-value of ~0.06 or less.
- HgCdTe is special in that it has an excess noise factor of 1 when injecting electrons.
- Some companies or researchers choose to use a slightly different notation:

 $M^{x}=kM+(1-k)(2-M^{-1})$

 In these cases, the quoted "excess noise factor" (or excess noise figure) is denoted as x instead of F.



Performance Metrics (GM-APD)

Avalanche Initiation Probability (→ PDE)

The probability that a carrier in a sufficiently high electric field will initiate a selfsustaining avalanche is the "avalanche initiation probability."

It is a function of:

- the strength of the electric field (specifically, how much the region is biased past the breakdown voltage, or the overbias)
- the location of the carrier in the electric field region
- the electron / hole ionization rates (depending on how the device is set up, either electrons or holes will initiate the avalanche)

Along with external and internal QE, the avalanche initiation probability defines the Photon Detection Efficiency (PDE), which is a standard metric for characterizing APDs.



Performance Metrics (GM-APD)

Dark Count Rate (DCR)

Dark count rate in Geiger mode APDs is similar to that in linear mode APDs. Thermally-generated dark carriers are multiplied and are indistinguishable from signal. The dark count rate must be characterized to account for the expected number of dark counts when taking data.



Performance Metrics (GM-APD)

Afterpulsing

Imperfections in the silicon lattice cause "trap" sites, which generate trap states (energy levels) into which carriers can "fall."

Traps have characteristic de-trapping lifetimes depending on type (lattice mismatch, surface termination / interface imperfections, foreign materials like metals, etc.). Once a carrier enters a trap state, the length of time it remains in that state has an exponential decay probability function.

Traps are filled during the avalanche surge, and when the carriers are released later on in the detection cycle, they can cause another avalanche at random.



Performance Metrics (GM-APDs)

Afterpulsing (self re-triggering)

A second type of afterpulsing arises from the emission of photons by carriers participating in an avalanche.

A silicon APD avalanche has an emission spectrum roughly equivalent to a black body at 4200 K, which has a substantial contribution at wavelengths that can be absorbed by silicon.

These photons are absorbed elsewhere inside the detector (perhaps in a neighboring pixel or in the same pixel), and can cause afterpulsing and optical crosstalk.



Quenching

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Passive Quenching:

A passive quenching circuit uses an in-series resistor to limit the current flow and decrease the voltage across the APD.

- Before an avalanche, the circuit is biased so that the reverse bias voltage across the APD is above breakdown. No current flows, so the supply voltage level drops completely across the APD.
- 2. When an avalanche occurs, high current is drawn, though limited by the resistor. As the current drawn by the APD increases, the voltage drop across the resistor increases. This, in turn, decreases the voltage across the APD, eventually decreasing it below breakdown and the avalanche ceases.
- 3. Current continues to flow through the resistor to charge the APD (which acts as a capacitor under reverse bias when not avalanching) until the voltage across the APD returns to the supply voltage level.



Quenching

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Active Quenching:

An active quench circuit senses the increase in current through the circuit (usually by a drop in voltage across an in-series resistor) and actively pins the voltage across the APD to below breakdown voltage for a set period.

- 1. The avalanche is sensed by the active quenching circuit and the APD voltage is immediately decreased below breakdown. This quenching voltage is held for a predetermined amount of time to evacuate all the avalanching carriers from the multiplication region.
- 2. At the end of the quench time, the voltage is either allowed to increase via RC recharge through a resistor (passive recharge), or it is actively set above breakdown voltage (active recharge).



Integrated Operation (GM-APD)

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Detection / Clocking Cycle (passive quenching)



Integrated Operation (GM-APD)

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Detection / Clocking Cycle (active quenching)





Integrated Operation (GM-APD)

Dead Time

Dead time refers to the period of the clocking cycle when the detector cannot register a photon event.

Generally, dead time includes quench time and the active reset time. Increased dead time reduces SNR and requires longer integration times to offset the loss of signal for constant sensitivity.



LIDAR

LIDAR (Light Detection and Ranging) uses the time of flight from a transmitter laser to calculate the distance to and placement of distant objects in three-dimensional space, most commonly in airborne applications.



3D LIDAR Image of a Portion of the RIT Campus





Imaging

Both linear mode and Geiger mode APDs have strengths and weaknesses when it comes to imaging.

- GM-APDs can be implemented with zero read noise (often the limiting factor in state-ofthe-art imaging systems) when designed with an in-pixel counter. This makes the entire detection process digital.
- GM-APDs cannot distinguish between one or more photons arriving simultaneously, while that information is maintained in a LM-APD.
- LM-APDs have effectively no dead time, which this is a limiting factor in GM-APD performance.
- LM-APDs must be reset periodically to remove the charge accumulated by leakage current or the detector quickly becomes saturated, while this is not a problem with GM-APDs.

0 and 10 electrons read noise for images with simulated shot noise





Imaging (single photon counting)



Figures Courtesy of Don Hall (University of Hawaii)



Imaging (GM-APD)



This image taken with a prototype GM-APD imager by researchers at Lincoln Laboratory (MIT).

Image courtesy of Don Figer.



Currently Available Hamamatsu S11519 Series



General ratings / absolute maximum ratings

Type no.				Absolute maximum ratings		
	Window material*1	Package	Active area size*2	Operating temperature Topr	Storage temperature Tstg	
			(mm)	(°C)	(°C)	
S11519-10	K	TO-5	φ1.0	-20 to +85	-55 to +125	
S11519-30	К	TO-8	φ3.0	-20 10 +65		

*1: K=borosilicate glass

*2: Area in which a typical gain can be obtained

Electrical and optical characteristics (Typ. Ta=25 °C, unless otherwise noted)

Type no.	Spectral response range λ	Peak sensitivity wavelength* ³ λp	Breakdown voltage VBR ID=100 µA		Temp. coefficient of VBR	Dark current* ³ ID		Terminal capacitance ^{*3} Ct	Cut-off frequency*3 fc	Excess noise figure* ³	Gain M
			Тур.	Max.	10-100 pA	Тур.	Max.		NL-30 32	λ-800 nm	~−090 mm
	(nm)	(nm)	(V)	(V)	(V/°C)	(nA)	(nA)	(pF)	(MHz)	7090 mm	
S11519-10	600 to 1150	960	350	500	1.7	3	30	2.0	400	0.3	100
S11519-30				500		9	90	12.0	230		100

*3: Values measured at a gain listed in the characteristics table

 $M^{x}=kM+(1-k)(2-M^{-1})$



Spectral response

Currently Available

Hamamatsu S11519 Series





Dark current vs. reverse voltage





Reverse voltage (V)

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